

octahedral arrangement of oxygen atoms), the O–Na–O angles vary from 55 to 140° with the smallest angle involving oxygen atoms which share an edge with the sulfate tetrahedron. The next smallest angle (66°) involves the oxygen atoms forming the smallest inter-sulfate separation mentioned above; hence, this edge is shared by three identical, but rotated, polyhedra. However, such an octahedral environment appears to be too restrictive for the coordination about the sodium ion. By considering the next two closest oxygen atoms [O(6') and O(4')] in the coordination sphere, a distorted Archimedean antiprism (Fig. 2) is obtained. In this idealized polyhedron, the O(4) type oxygen atoms are displaced toward the sodium ion from the top plane of that required for a regular antiprism. The calcium ion in $\text{CaNa}_2(\text{SO}_4)_2$ (Cocco, Corazza & Sabelli, 1965) is also situated in a distorted antiprism consisting of six closer (2.36–2.47 Å) and two more distant (2.74 Å) oxygen atoms.

Lithium sodium sulfate belongs to a class of compounds of the form $M'M''\text{XO}_4$ which are all trigonal or higher in symmetry, but belong to different space groups, and which have lattice constants related by simple multiples. Among these compounds are: (1) LiNaSO_4 in space group $P31c$; (2) KLiSO_4 and KAlSiO_4 in space group $P6_32$ (Claringbull & Bannister (1948) showed these to be related to hexagonal tridymite, nepheline and kaliophilite); (3) $\text{NaK}_3(\text{SO}_4)_2$ and NaKSO_4 in space group $P\bar{3}m1$ (Bellanca, 1943); (4) anhydrous alums such as $\text{KAl}(\text{SO}_4)_2$ in space group

$P321$ (Vegard & Maurstad, 1929); and (5) silicates such as NaAlSiO_4 and KLaSiO_4 in space group $P6_3$ (subcell to KLiSO_4 ; Eitel, Herlinger & Trömel, 1930). With the exception of the last set of silicates, for which the detailed crystal structure has not been proposed, all these compounds have sulfate or silicate groups situated on the threefold axes. The relative locations of the XO_4 groups appear to be governed by the manner in which the metal ions form their coordination polyhedra.

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The Crystal Structure of L-Leucine Hydrobromide*

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The crystal structure of L-leucine hydrobromide has been determined. The orthorhombic unit cell has the dimensions $a = 7.29$, $b = 24.51$, $c = 5.54$ Å and the space group is $P2_12_12_1$. There are four molecules of $\text{C}_6\text{H}_{14}\text{NO}_2^+ \cdot \text{Br}^-$ in the cell. The trial structure, obtained by a combination of the 'heavy atom' method and the method of structure invariants applied to the c -axis projection, has been refined by three-dimensional least-squares method to an R value of 0.11. All available hydrogen atoms are involved in hydrogen bonding.

Introduction

The investigation of the crystal structure of leucine was taken up as part of a program of research work in this

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department on the structure analysis of amino acids and proteins. Leucine is among the most important amino acids essential for the growth and maintenance of living organisms, and, while the crystal structures of most of the amino acids have been determined, the structure analysis of leucine has not so far been reported except as an N -terminal residue in the tripeptide leucylprolylglycine (Leung & Marsh, 1958). Knowledge of the stereochemistry of leucine as a structural unit in proteins would be of great value and consequently the structure analysis was taken up.

Experimental

Crystals of leucine hydrobromide were obtained by dissolving L-leucine in 30% hydrobromic acid and allowing the solution to evaporate. The crystals grow as thin plates with principal faces (010) and elongated parallel to the *c* axis. They are very soft, highly hygroscopic, and decompose quickly on exposure to air; hence the freshly prepared crystals were immediately sealed in Lindemann-glass capillary tubes. X-ray photographs taken with equi-inclination Weissenberg as well as precession methods established the crystallographic data summarized in Table 1. The space group absences were of the type $h00$, h odd; $0k0$, k odd; and $00l$, l odd. Three-dimensional intensity data for Cu $K\alpha$ radiation were obtained from multiple-film equi-inclination Weissenberg photographs, with the crystal rotated about the *c* axis ($l=0$ to 4). Of the 1098 reflexions which could have been recorded in the various layers, only 593 were observed. The intensities were estimated visually by using calibrated filmstrips and were corrected for Lorentz and polarization factors by means of the chart due to Cochran (1948). Except for a few very low-angle reflexions, no spot-size correction was applied. The linear absorption coefficient of the crystal for Cu $K\alpha$ radiation is 59.8 cm^{-1} . The crystal used had a cross-section of $0.01 \times 0.02 \text{ cm}^2$. Thus the maximum error in the relative values of the structure factors because of absorption is about 30%, although the average error will be considerably smaller than this. However no corrections for absorption were made.

Table 1. *Crystal data for L-leucine hydrobromide*

$\text{C}_6\text{H}_{13}\text{NO}_2 \cdot \text{HBr}$ Orthorhombic	M.W. 212.1 Space group $P2_12_12_1$
$a = 7.29 \text{ \AA}$	$Z = 4$
$b = 24.51$	$\rho_o = 1.423 \text{ g.cm}^{-3}$
$c = 5.54$	$\rho_c = 1.435 \text{ g.cm}^{-3}$
$\lambda = \text{Cu } K\alpha (1.5418 \text{ \AA})$	$F_{000} = 436$

Photographs about the *a* axis could not be taken since the crystals disintegrated on cutting. Consequently, a crystal was mounted along the *c* axis on a precession camera and the $0kl$ zone was recorded with Cu $K\alpha$ radiation; this photograph had reflexions up to $l=3$. The intensities were estimated visually, corrected for Lorentz and polarization factors by the tables given by Nordman (1959), and used to correlate intensities from the Weissenberg photographs for $l=0$ to 3. Approximate scale and temperature factors ($B=5.0 \text{ \AA}^2$) were derived by the method of Wilson (1942); final scaling, however, was achieved by comparing $\Sigma |F_o|$ with $\Sigma |F_c|$ for each layer. Atomic scattering factors of Berghuis, Haanappel, Potters, Loopstra, MacGillavry & Veenendaal (1955) for carbon, nitrogen, and oxygen and of Thomas & Umeda (1957) for bromine, were used in projection calculations.

Derivation of the trial structure

A *c*-axis Patterson projection gave the approximate coordinates $x=0.250$, $y=0.073$ for the bromine atom, and a bromine-phased electron-density projection was calculated. However, with $x=0.25$, the bromine atom does not contribute to $hk0$ reflexions having $h+k$ odd, and the omission of these terms from the Fourier synthesis resulted in spurious mirror planes at $x=\frac{1}{4}$ and $\frac{3}{4}$. Since most of the light-atom peaks lay near these planes, it was difficult to choose between a peak and its spurious mirror-related mate, and continued attempts to derive a satisfactory structure were unsuccessful. The structure was eventually solved by the application of the method of structure invariants (Woolfson, 1957).

Application of the method of structure invariants

It has been shown by various workers (*e.g.*, Sayre, 1952; Cochran, 1952; Zachariasen, 1952; Hughes, 1953) that the signs of the largest unitary structure factors [$U(\mathbf{h})$, $U(\mathbf{h}')$ and $U(\mathbf{h}+\mathbf{h}')$] are most probably related by the equation

$$s(\mathbf{h}')s(\mathbf{h}+\mathbf{h}') \sim s(\mathbf{h}). \quad (1)$$

The method of structure invariants consists in finding a series of relations of the type (1) in which the $s(\mathbf{h})$'s now represent the signs of the structure-invariant reflexions, *i.e.* reflexions whose signs are independent of the choice of origin. The structure-variant reflexions are related in pairs with a structure invariant, so that if the signs of the latter are known, these relationships lead to a series of sign linkages between the former.

The centrosymmetric (001) projection for leucine hydrobromide has the plane-group symmetry pgg and the signs of the structure-invariant reflexions are known (since the bromine atom contributes to all reflexions with $h+k=2n$). The unitary structure factors for the $hk0$ reflexions were calculated by using the expression

$$|U(\mathbf{h})| = \frac{|F_o(\mathbf{h})|}{\sum_i f_i} \left[\frac{\sum f_i^2}{\langle I \rangle} \right]^{\frac{1}{2}},$$

where $|F_o(\mathbf{h})|$ is the relative structure amplitude of the $hk0$ reflexion. In applying the above expression, the procedure adopted was to subtract the bromine contribution from the scaled F_o values for reflexions of the type with $h+k=2n$, thus assuming that the signs of these reflexions were the same as those given by bromine. However, only those reflexions for which $|F_{\text{Br}}(\mathbf{h})|/|F_o(\mathbf{h})|$ was greater than 0.5 were included. In this way the problem is reduced to that of solving a light atom structure with the difference that the signs of half of the reflexions are already known.

Those terms for which $|U(\mathbf{h})| > 0.20$ were then selected and the method of Woolfson was applied to obtain the signs of 14 $hk0$ reflexions with $(h+k)$ odd.*

* In the final analysis, the signs of three of these reflections were found to be wrongly determined.

These terms were then included in a Fourier summation for the c -axis projection, which is shown in Fig. 1. The spurious mirror symmetry had completely broken down and the map could be very easily interpreted in terms of the known chemistry of the molecule. The resulting trial structure has been plotted in Fig. 1; it gave an R index of 0.25 for the $hk0$ reflexions. A difference-Fourier synthesis drawn at this stage indicated large shifts for a number of atoms and pronounced anisotropy of thermal motion for the bromine atom. The anisotropy was, however, neglected but the shifts in the positional parameters were estimated and a fresh set of $hk0$ structure factors was calculated. A few cycles of Fourier refinement brought the R index for these reflexions down to 0.22.

The z coordinate of the bromine atom was obtained from a Patterson projection on (100), calculated from the relatively few $0kl$ reflexions available on the upper-level Weissenberg photographs about c . The z coordinates of the lighter atoms were then deduced by fitting a model of the molecule to the c -axis and a -axis electron-density projections (Figs. 1 and 2).

Three-dimensional least-squares refinement

Preliminary structure-factor least-squares calculations were carried out on an Elliott 803 computer and, later, on a CDC-3600, in each case using a diagonal least-squares program written by the author specifically for the space group $P2_12_12_1$. The function minimized was $\sum w(|F_o| - |F_c|)^2$, with unit weights assigned at the start. The scattering factors were calculated by an analytical

approximation due to Vand, Eiland & Pepinsky (1957) and with the constants given by Forsyth & Wells (1959).

The first three-dimensional calculation led to an R index of 0.26 (throughout, $R = \sum ||F_o| - |F_c|| / \sum |F_o|$); after several cycles of refinement, in which the individual isotropic temperature factors as well as the atomic coordinates were adjusted, convergence was reached at $R = 0.18$. During this refinement the temperature factor for C(3) [C(3) is the carbon atom of the methyl group attached to C(2); C(1) is that of the other methyl group (Fig. 4)] reached the unrealistically large value of 10.2 \AA^2 , raising suspicion that this atom might have been misplaced. A three-dimensional difference map was therefore calculated, the contributions of C(3)

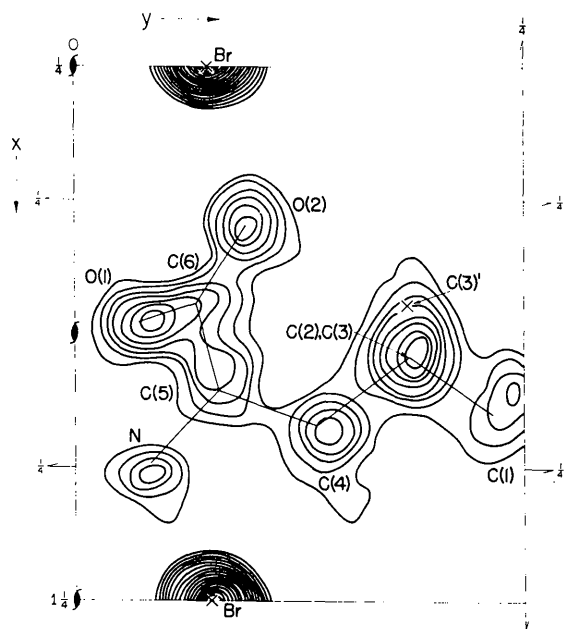


Fig. 1. c -Axis bromine-phased electron-density projection, with signs of some of the reflexions with $h+k$ odd determined by the method of structure invariants. Contours are at arbitrary equal intervals.

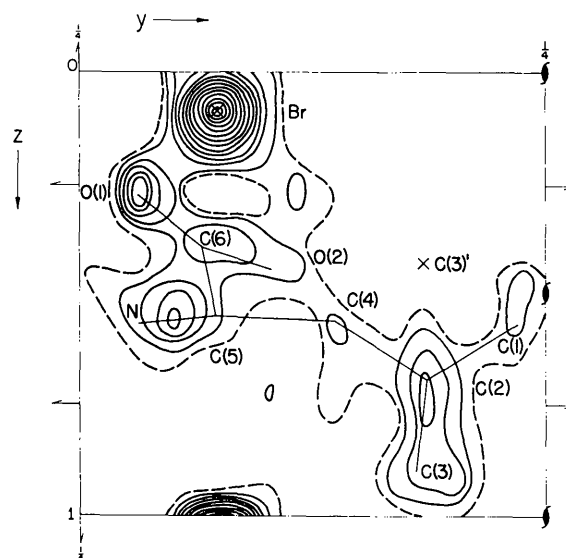


Fig. 2. a -Axis electron-density projection, phased on the contribution from the bromine atom.

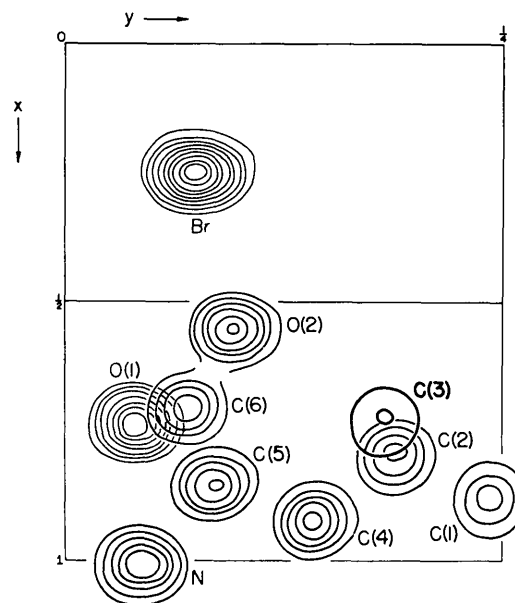


Fig. 3. Final electron density projected down the a axis.

having been omitted from the values of F_c . This map showed a peak of approximately one-half the expected height at the position assumed for C(3), and a second peak of about the same height approximately 1.4 Å from C(2) and roughly along the remaining tetrahedral direction (presumably occupied by the hydrogen atom). Although this second peak was very broad, it suggested that C(3) might be disordered between two positions corresponding to a random interchange between the hydrogen and C(3) in different molecules. This peak is indicated by C(3)' in Figs. 1 and 2 (the crosses).

Two atoms with half the carbon scattering power were therefore assumed at C(3) and C(3)' and the least-squares refinement was continued. It was found that while the thermal parameter for C(3) was quite normal, that for C(3)' increased to 22.0 Å², although the R value itself converged to 0.16.

It was therefore felt that the abnormal value of B for C(3)' probably indicated the absence of any atom there, and that the peak for C(3)' mentioned earlier might be spurious. In addition, the short interatomic distance (~ 2.7 Å) between C(3)' of one molecule and C(1) of a molecule related by the twofold screw axis parallel to the a axis makes it improbable that there would be any disorder at all. It appears that the large temperature factor of C(3) merely reflects a large thermal motion due to loose packing in this region of the structure.

It was therefore decided to omit C(3)' and to continue the refinement with full strength at C(3). The atomic coordinates were the same as those used at the beginning of the cycles of refinement which included C(3)'. With a modified Hughes (1941) weighting scheme convergence was soon reached at $R=0.16$.

Refinement was concluded at the California Institute of Technology, using the CRYRM system (Duchamp, 1964) on an IBM 7094 computer. The quantity minimized was $\sum w(F_o^2 - F_c^2)^2$ with weights taken according to the expressions

$$w = 1/F_o^2 \quad F_o > 40$$

$$w = 1/40F_o \quad F_o < 40.$$

Unobserved reflexions were included only when the calculated structure factors exceeded the visual threshold.

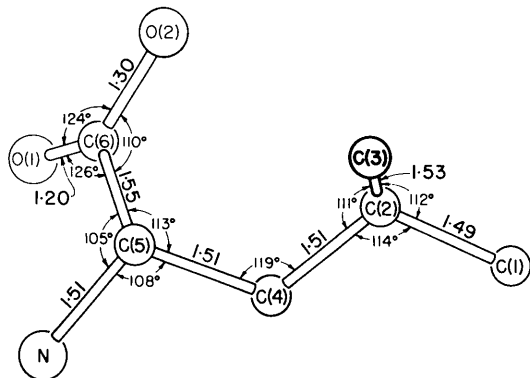


Fig. 4. Bond distances and bond angles observed in the molecule of leucine. E.s.d.'s are 0.03–0.04 Å and 4° respectively.

Three cycles of full-matrix refinement with individual isotropic temperature factors for all the atoms resulted only in a slight decrease in R from 0.211 to 0.198 (the higher index being due to the inclusion of many unobserved reflexions). Six cycles of further refinement in which the bromine atom was allowed to become anisotropic brought the R index down to 0.139.

A three-dimensional Fourier synthesis calculated at this stage failed to show any strength near the C(3)' position, thus confirming the earlier observation regarding its spurious nature. Difference Fourier syntheses were then calculated for sections likely to contain the hydrogen atoms. Except for the hydrogen atoms attached to C(1) and N, the positions of the hydrogen atoms were not apparent, and they were therefore fixed purely from spatial considerations; no attempt was made to determine the positions of the hydrogen atoms attached to C(3) in view of the latter's large thermal motion. These hydrogen atoms were included in the subsequent structure factor calculations with their thermal parameters assumed to be the same as those of the atoms to which they are covalently bonded. A few cycles of full-matrix anisotropic refinement were then carried out for all the light atoms ($R=0.122$); however, the thermal ellipsoids for most of the atoms did not make physical sense. Consequently the anisotropic refinement for the light atoms was ignored and all of them were refined isotropically again until no shift was greater than 10% of its e.s.d. The final R index for all the reflexions was 0.113. Fig. 3 shows the final electron density projected down the a -axis.

Table 2 lists the F_o 's and F_c 's, while Tables 3 and 4 give the final parameters for all the atoms.

Discussion

In spite of the large standard deviations (0.04 Å in C–C bonds and 0.03 Å in C–N and C–O bonds; $\sim 4^\circ$ in bond angles), due mainly to lack of sufficient data and the predominant scattering power of bromine, all the bond distances and angles are quite reasonable (Fig. 4). The carboxyl group and the α -carbon atom [C(5)] are coplanar within limits of accuracy of the data. The equation of the mean plane passing through them, in terms of the direction cosines and the origin-to-plane distance, is given by

$$-0.320X - 0.780Y + 0.538Z = 1.705 \text{ \AA}.$$

The deviations of the individual atoms from this plane are C(5) = -0.01 Å, C(6) = 0.04 , O(1) = -0.02 , and O(2) = -0.01 . The nitrogen atom is 0.24 Å away from this plane. The conformation about the various bonds is shown in Fig. 5. It is seen that all the bonds assume a staggered conformation.

Intermolecular arrangement

The intermolecular packing arrangement, viewed down the c and a axes, is shown in Figs. 6 and 7. The hydrogen bonds are indicated by dashed lines. Table 5

lists the hydrogen bond distances and angles. Fig.8 shows a stereoscopic view of the molecule, together with the hydrogen bonding.

There are four hydrogen atoms available for hydrogen-bond formation — three associated with the amino

nitrogen atom and one with the carboxyl oxygen atom O(2). All of them are utilized in hydrogen bonding. The nitrogen atom is at hydrogen-bonding distance from four atoms — three bromide ions and an oxygen atom [Br(II), Br(III), Br(IV), and O(1)(IV) in Table 5].

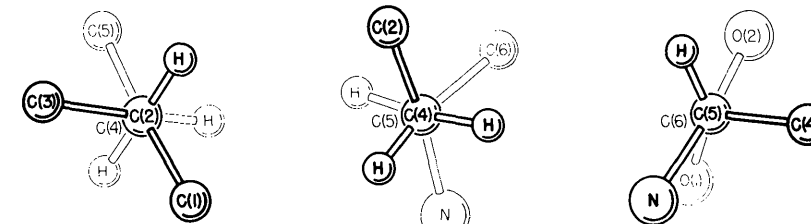


Fig. 5. Conformation about the various bonds in the molecule of leucine.

Table 2. Observed and calculated structure factors
Values given are 10 |Fo|, 10 |Fc| and α .

Table with 4 columns: h k l, 10 |Fo|, 10 |Fc|, and alpha. The table contains a large number of rows of numerical data representing structure factors for various Miller indices.

Table 3. *The final parameters and their estimated standard deviations*For the bromine atom the temperature factor is of the form $T = \exp[-(b_{11}h^2 + b_{22}k^2 + b_{33}l^2 + b_{12}hk + b_{13}hl + b_{23}kl)]$.

	x	y	z	B
Br	0.2478 (4)	0.0741 (1)	0.0796 (5)	*
O(1)	0.7385 (26)	0.0396 (5)	0.2546 (28)	5.0 (3) Å ²
O(2)	0.5592 (26)	0.0943 (7)	0.4657 (35)	6.3 (4)
N	1.0130 (25)	0.0439 (7)	0.5837 (38)	4.5 (4)
C(1)	0.8790 (51)	0.2421 (12)	0.6065 (55)	7.5 (7)
C(2)	0.7999 (35)	0.1871 (10)	0.6592 (52)	5.7 (6)
C(3)	0.7312 (57)	0.1826 (16)	0.9188 (73)	10.7 (11)
C(4)	0.9246 (32)	0.1403 (9)	0.5965 (50)	5.0 (5)
C(5)	0.8546 (28)	0.0828 (8)	0.6242 (40)	3.8 (4)
C(6)	0.7103 (25)	0.0674 (8)	0.4297 (43)	4.1 (4)

For Br $b_{11} = 0.0190$ (4)
 $b_{22} = 0.0034$ (1)
 $b_{33} = 0.0452$ (13)
 $b_{12} = 0.0007$ (4)
 $b_{13} = 0.0117$ (17)
 $b_{23} = 0.0006$ (4)

Table 4. *Coordinates (fractional) of the hydrogen atoms*H(*ij*) indicates the *j*th hydrogen attached to the *i*th atom.

	x	y	z
H(11)	0.84	0.27	0.72
H(12)	1.01	0.24	0.61
H(13)	0.88	0.25	0.44
H(21)	0.71	0.19	0.54
H(41)	1.04	0.14	0.70
H(42)	0.97	0.14	0.42
H(51)	0.77	0.08	0.80
H(71)	1.11	0.05	0.69
H(72)	1.07	0.05	0.43
H(73)	0.98	0.01	0.59
H(81)	0.44	0.11	0.51

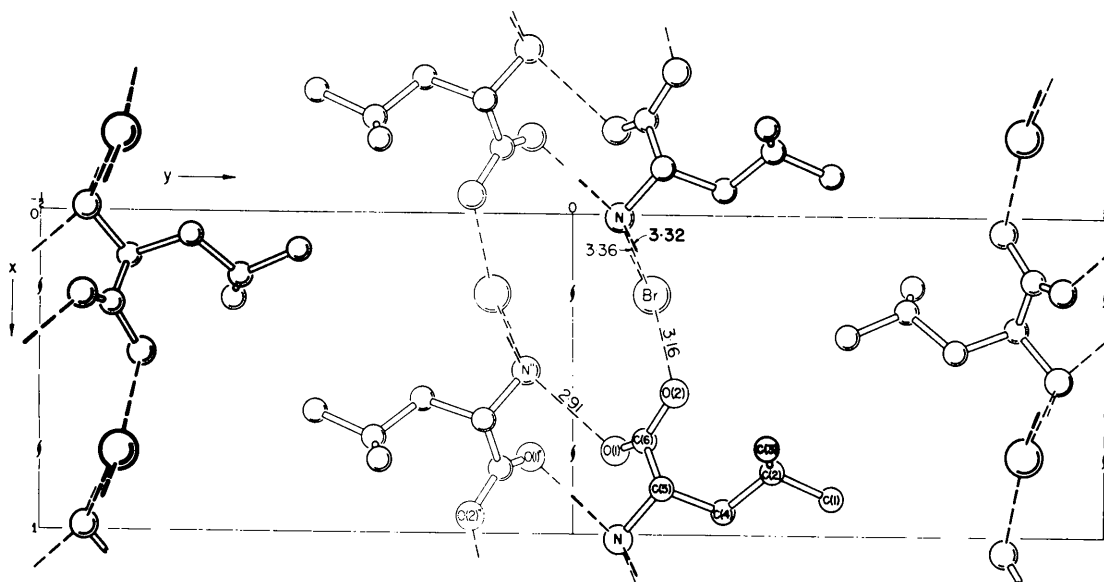
Table 5 (cont.)

N(1)···O(1)(IV)	2.91 Å
O(2)(I)···Br(I)	3.16
N(I)···Br(IV)	3.38
C(5)(I)-N(I)-Br(II)	112°
C(5)(I)-N(I)-Br(III)	98
C(5)(I)-N(I)-Br(IV)	159
C(5)(I)-N(I)-O(1)(IV)	85
Br(II)-N(I)-Br(III)	112
Br(II)-N(I)-Br(IV)	85
Br(III)-N(I)-Br(IV)	86
O(1)(IV)-N(I)-Br(II)	138
O(1)(IV)-N(I)-Br(III)	102
O(1)(IV)-N(I)-Br(IV)	74
C(6)(I)-O(2)(I)-Br(I)	115
C(6)(IV)-O(1)(IV)-N(I)	124

Table 5. *Hydrogen bond distances and angles*

I	x, y, z	
II	$x+1, y, z$	
III	$x+1, y, z+1$	
IV	$-x+\frac{3}{2}, -y, z+\frac{1}{2}$	
N(1)···Br(II)	3.36 Å	
N(I)···Br(III)	3.32	

Two of the hydrogen bonds [to Br(II) and Br(III)] are nearly in the tetrahedral direction. Regarding the third bond, the angles C(5)-N-Br(IV) and C(5)-N-O(1)(IV) (159 and 85° respectively) together with the comparatively large deviations from the tetrahedral angle

Fig. 6. Intermolecular packing viewed down the *c* axis.

of Br(II)-N-Br(IV) and Br(III)-N-Br(IV) as compared with Br(II)-N-O(1)(IV) and Br(III)-N-O(1)(IV) (Table 5) seem to indicate a hydrogen bonding to the oxygen atom only, and not to the bromine atom. These considerations also rule out the possibility of bifurcation of the hydrogen bond between the oxygen and the bromine. The N...Br(IV) distance of 3.38 Å would then have to be treated as a short contact rather than

a hydrogen bond. Short contacts occurring under similar circumstances have been reported (Donohue, 1952). In these cases, after all the hydrogen atoms on the $-\text{NH}_3^+$ group have been used up in hydrogen-bond formation, a short contact occurs with a chloride ion which is generally situated (at hydrogen bond distance) approximately on the extension of the X-N covalent bond, so that the angle X-N...Cl⁻ is near 180°.

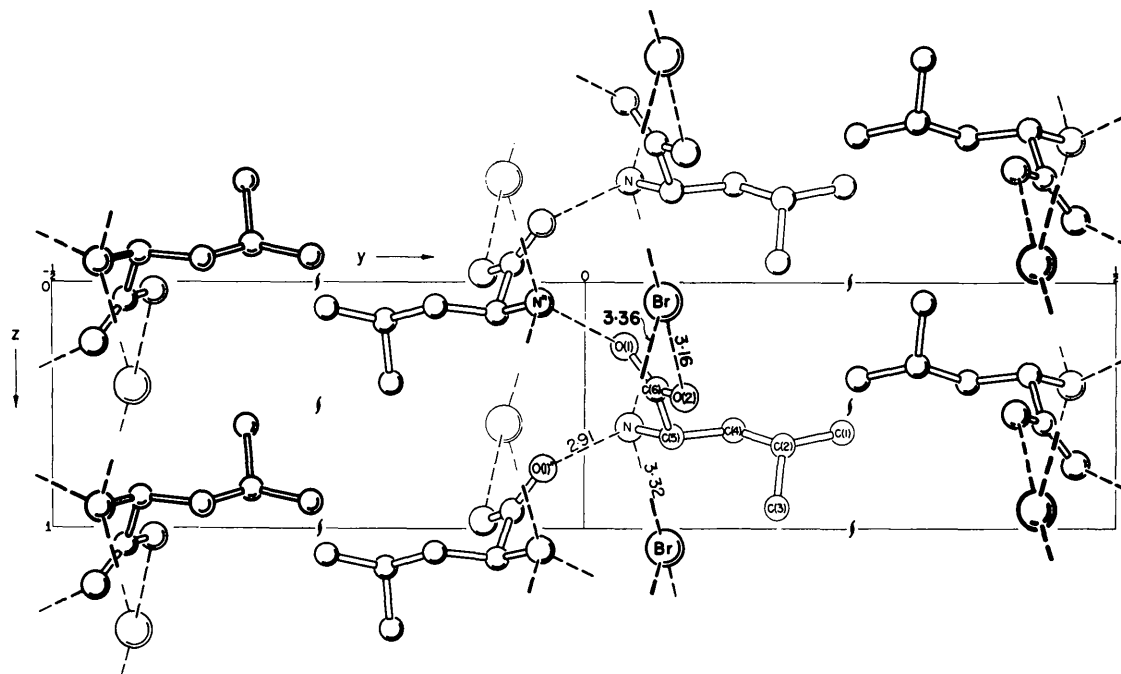


Fig. 7. Intermolecular packing viewed down the *a* axis.

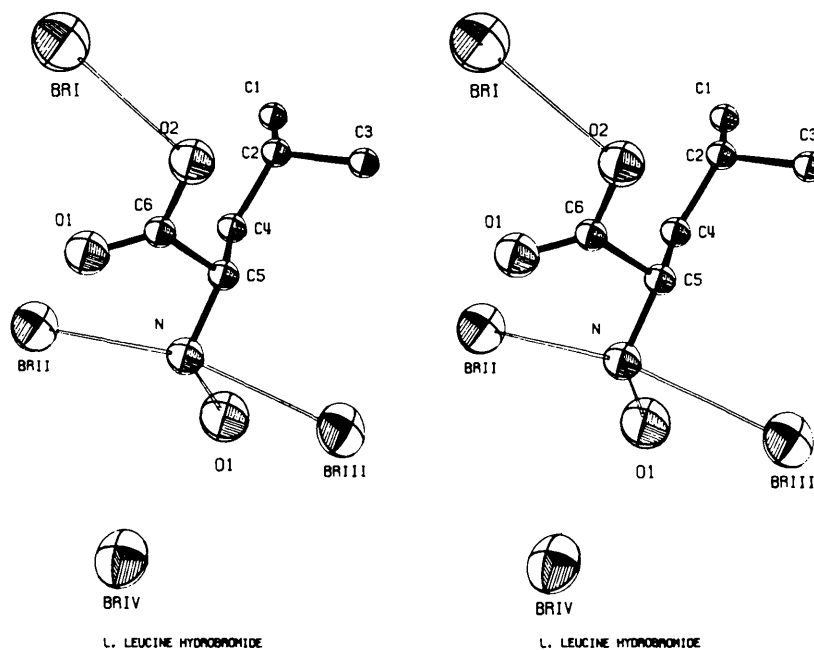


Fig. 8. Stereoscopic view of the molecule.

It is interesting to note that the hydrogen bond network forms sheets of molecules perpendicular to the *b* axis. The sheets are held together only by van der Waals forces. Intermolecular contacts less than 4 Å are listed in Table 6.

Table 6. Intermolecular contacts less than 4 Å

Roman numerals indicate the relation of the second atom to the first at (*x*, *y*, *z*)

Atom 1	Atom 2	Distance
	I $\frac{1}{2}+x, \frac{1}{2}-y, 1-z$	
	II $\frac{3}{2}-x, -y, \frac{1}{2}+z$	
	III $x, y, z+1$	
	IV $x+1, y, z+1$	
	V $x-1, y, z$	
C(1)	C(1)(I)	3.85 Å
C(1)	C(2)(I)	3.82
N	C(6)(II)	3.71
C(5)	O(1)(II)	3.16
C(6)	O(1)(II)	3.20
N	O(1)(II)	2.91
O(2)	O(1)(II)	3.94
O(1)	O(1)(II)	3.39
N	Br(II)	3.38
C(3)	O(2)(III)	3.93
C(3)	O(1)(III)	3.97
C(5)	O(1)(III)	3.75
C(4)	Br(IV)	3.92
C(5)	Br(IV)	3.82
N	Br(IV)	3.32
Br	N(V)	3.36
Br	O(1)(V)	3.93

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